

# Magnetic Susceptibilities of Cubic Mixed Europium Oxides

A. GRILL AND M. SCHIEBER

*Department of Physics, Hebrew University of Jerusalem, Jerusalem, Israel*

(Received 17 October 1969)

The magnetic susceptibility  $\chi$  of polycrystalline cubic solid solutions of  $\text{Eu}_2\text{O}_3$  with  $\text{Y}_2\text{O}_3$  and  $\text{Gd}_2\text{O}_3$  have been measured. The  $\chi$  of  $\text{Eu}^{3+}$  in  $(\text{Eu}_x\text{Y}_{1-x})_2\text{O}_3$  increases with decreasing concentration of  $\text{Eu}^{3+}$  ions, although the anisotropic exchange interaction would be expected to decrease at the same time. The  $\chi$  of the smaller unit-cell compound,  $(\text{Eu}_{0.1}\text{Lu}_{0.9})_2\text{O}_3$ , yields an even higher value for the  $\chi$  of  $\text{Eu}^{3+}$  than do the Eu-Y mixed oxides, indicating a strong dependence of the crystal-field parameters on the unit-cell dimension. In  $(\text{Eu}_x\text{Gd}_{1-x})_2\text{O}_3$  the  $\chi$  of  $\text{Eu}^{3+}$  has the same value as in  $\text{Eu}_2\text{O}_3$  for all values of  $x$  while the  $\chi$  of  $\text{Gd}^{3+}$  is given by  $\chi = C/[T + (1-x)\theta]$ .

## INTRODUCTION

THE increased magnetic susceptibility  $\chi$  of the  $\text{Eu}^{3+}$  ion in  $\text{Eu}_2\text{O}_3$  as compared to its free-ion value<sup>1</sup> was interpreted recently by Huang and Van Vleck<sup>2</sup> in terms of the combined action of the crystalline field and anisotropic exchange. It has been assumed that the  $\chi$  of the  $\text{Eu}^{3+}$  ion diluted by diamagnetic ions of  $\text{Y}^{3+}$  or  $\text{Lu}^{3+}$  should decrease at low concentrations of  $\text{Eu}^{3+}$  because of the decrease of the anisotropic exchange interaction between the  $\text{Eu}^{3+}$  ions. The present work shows that the  $\chi$  of  $\text{Eu}^{3+}$  is increased at low concentrations of  $\text{Eu}^{3+}$  in the solid solutions of Eu-Y and Eu-Lu oxides, and remains almost unchanged at higher concentrations of  $\text{Eu}^{3+}$  in the Eu-Y oxides and at all concentrations of  $\text{Eu}^{3+}$  in the Eu-Gd oxides.

Huang and Van Vleck<sup>2</sup> were handicapped by insufficient experimental data on the dependence of the  $\chi$  of  $\text{Eu}^{3+}$  on concentration, and it is to remedy this deficiency that the present work was undertaken at their suggestion.

The cubic rare-earth oxides crystallize in the bixbyite structure. There are two different sites for the metal ions; 75% occupy the asymmetric  $C_2$  site, and the remainder occupy the more symmetric  $S$  site. The present work can be interpreted as indicating that the  $\text{Eu}^{3+}$  ions prefer the  $C_2$  sites for low concentrations of  $\text{Eu}^{3+}$  in the mixed cubic oxides of Y and Lu.

## EXPERIMENTAL PROCEDURE

The solid solutions  $(\text{Eu}_x\text{Y}_{1-x})_2\text{O}_3$  and  $(\text{Eu}_x\text{Gd}_{1-x})_2\text{O}_3$  with  $x=0.2, 0.4, 0.6, 0.8$ , as well as  $(\text{Eu}_{0.1}\text{Lu}_{0.9})_2\text{O}_3$ , were prepared by mixing the individual oxides in the same molar ratios as in the final product, and firing below the cubic-to-monoclinic phase-transition temperature of  $\text{Eu}_2\text{O}_3$  (1100°C). The specimens were resintered and refined until single cubic phases were shown in sharp powder diffraction patterns, taken with a Guinier x-ray camera. The unit-cell dimensions measured from the above patterns conform to Vegard's law, as shown in Fig. 1. Single-cubic-phase specimens  $(\text{Eu}_x\text{Lu}_{1-x})_2\text{O}_3$

with  $x$  greater than 0.1 were not obtained, because of the large difference between the unit cells of the two pure oxides (10.87 and 10.39 Å).

The  $\chi$  of the specimens was measured with a null coil pendulum magnetometer<sup>3</sup> between 95 and 295°K on samples weighing about 200 mg.

## RESULTS AND DISCUSSION

The dependence of  $\chi$  on the concentration  $x$  in  $(\text{Eu}_x\text{Y}_{1-x})_2\text{O}_3$  and  $(\text{Eu}_{0.1}\text{Lu}_{0.9})_2\text{O}_3$  at 95°K is shown in Fig. 2. A similar variation of  $\chi$  with  $x$  is obtained also at the other temperatures. It can be seen that the  $\chi$  of  $\text{Eu}^{3+}$  is increased at low values of  $x$  and that the increase is more pronounced for the Eu-Lu than for the Eu-Y oxides.

The increased  $\chi$  of the  $\text{Eu}^{3+}$  ion at small concentrations can be explained partly by supposing that the crystal field is much stronger at the  $C_2$  than at the  $S_6$  site, and that at low concentrations the  $\text{Eu}^{3+}$  ions occupy only the  $C_2$  sites. The larger  $\chi$  for small values of  $x$  was expected to be diminished by the decrease of anisotropic exchange interaction between the  $\text{Eu}^{3+}$  ions, but this is contradicted by our experimental data. It is also possible to assume that the crystal field is largely re-

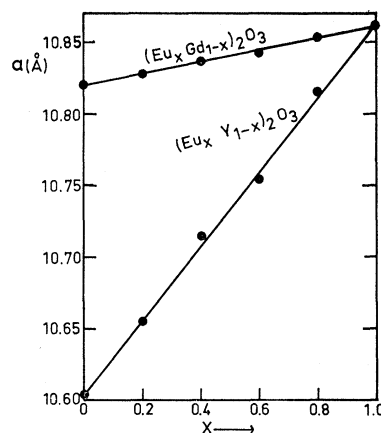


FIG. 1. Unit cell  $a$  versus concentration  $x$  of  $(\text{Eu}_x\text{Y}_{1-x})_2\text{O}_3$  and  $(\text{Eu}_x\text{Gd}_{1-x})_2\text{O}_3$ .

<sup>1</sup> R. M. W. Trapness and P. W. Selwood, *Nature* **169**, 840 (1952).

<sup>2</sup> N. L. Huang and J. H. Van Vleck, *J. Appl. Phys.* **40**, 1144 (1969).

<sup>3</sup> R. M. Bozorth, H. J. Williams, and D. E. Walsh, *Phys. Rev.* **103**, 572 (1956).

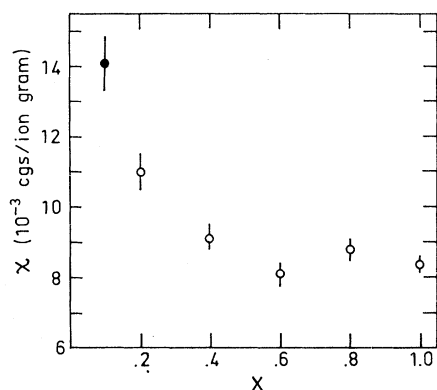


FIG. 2. Magnetic susceptibility  $\chi$  per ion gram of  $\text{Eu}^{3+}$  versus concentration  $x$  of  $(\text{Eu}_x\text{Y}_{1-x})_2\text{O}_3$  (circle) and of  $(\text{Eu}_{0.1}\text{Lu}_{0.9})_2\text{O}_3$  (dot) at  $95^\circ\text{K}$ .

sponsible for the increase, and that it is strongly dependent on the unit-cell dimensions. Indeed, this would agree with the finding of a larger  $\chi$  for the  $\text{Eu}^{3+}$  ion in  $(\text{Eu}_{0.1}\text{Lu}_{0.9})_2\text{O}_3$ , which has a smaller unit cell [ $a = 10.44 \text{ \AA}$  compared with  $a = 10.60 \text{ \AA}$  for  $(\text{Eu}_x\text{Y}_{1-x})_2\text{O}_3$  compounds]. The  $\chi$  of  $\text{Eu}^{3+}$  for  $x = 0.6$  is not markedly changed, presumably as a result of changes in the site occupation probabilities for  $\text{Eu}^{3+}$  and of increased unit-cell dimensions, compared to those in low- $x$  compounds. The dependence of the crystal-field parameters on the unit-cell dimensions is further confirmed in the  $(\text{Eu}_x\text{Gd}_{1-x})_2\text{O}_3$  compounds, where  $\chi$  of  $\text{Eu}^{3+}$  versus  $x$  is the same as in  $\text{Eu}_2\text{O}_3$  because of the small difference in the unit-cell dimensions ( $a = 10.82 \text{ \AA}$  and  $a = 10.87 \text{ \AA}$  for pure  $\text{Gd}_2\text{O}_3$  and  $\text{Eu}_2\text{O}_3$ , respectively, as compared with  $a = 10.61 \text{ \AA}$  and  $a = 10.39 \text{ \AA}$  for pure  $\text{Y}_2\text{O}_3$  and  $\text{Lu}_2\text{O}_3$ , respectively).

The temperature dependence of  $\chi$  for a concentrated ( $x = 0.8$ ) and a dilute ( $x = 0.2$ ) mixed Eu-Y oxide is shown in Fig. 3. The dependence of  $\chi$  on the separation  $E_{10}$  between the levels  $J = 1$  and  $J = 0$  is given ap-

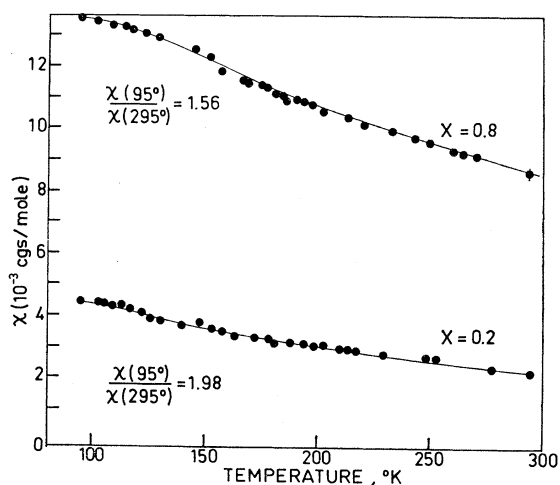


FIG. 3. Magnetic susceptibility  $\chi$  per mole of compound versus temperature  $T$  of  $(\text{Eu}_{0.2}\text{Y}_{0.8})_2\text{O}_3$  and  $(\text{Eu}_{0.8}\text{Y}_{0.2})_2\text{O}_3$ .

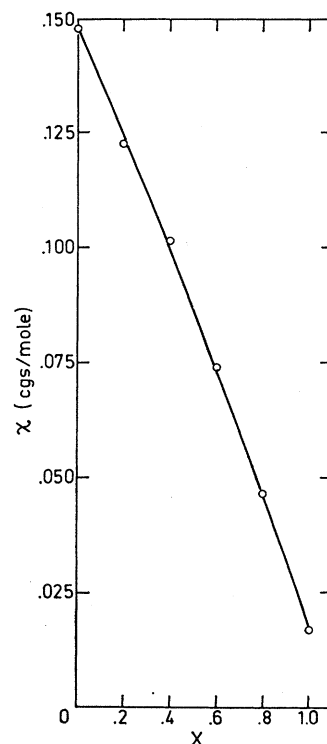


FIG. 4. Magnetic susceptibility  $\chi$  per mole of compound versus concentration  $x$  of  $(\text{Eu}_x\text{Gd}_{1-x})_2\text{O}_3$  at  $95^\circ\text{K}$ : circle represents experimental; line represents calculated with  $\chi = C/[T + (1-x)\theta]$ .

proximately by

$$\chi \sim [1 - \exp(-E_{10}/kT)]E_{10}^{-1}.$$

Our assumption is that  $E_{10}$  is larger at lower concentrations than at the higher ones. This means that  $\chi(95^\circ\text{K})/\chi(295^\circ\text{K})$  is larger at  $x = 0.2$  than at  $x = 0.8$ , as indeed is shown in Fig. 3.

The variation of  $\chi$  of  $(\text{Eu}_x\text{Gd}_{1-x})_2\text{O}_3$  with  $x$  at  $95^\circ\text{K}$  is shown in Fig. 4. The same variation of  $\chi$  with  $x$  is observed at all temperatures between 95 and  $295^\circ\text{K}$ . For  $x = 0$  the value of  $\chi$  is found to be given by  $C/(T + \theta)$ , with  $\theta = 11^\circ\text{K}$  rather than the  $18^\circ\text{K}$  reported previously by Velayos.<sup>4</sup> Reexamining Velayos's data, it can be seen that  $\theta = 11^\circ\text{K}$  seems to fit his own  $\chi$  at room and liquid-nitrogen temperatures better than  $\theta = 18^\circ\text{K}$ , which is obtained at high temperatures. The constant  $\theta$  represents the strength of the antiferromagnetic interaction between the  $\text{Gd}^{3+}$  ions and is proportional to the concentration  $(1-x)$  of  $\text{Gd}^{3+}$ . The  $\chi$  of  $\text{Gd}^{3+}$  is given therefore by  $\chi = C/[T + (1-x)\theta]$ , in good agreement with the experimental results as shown in Fig. 4.

#### ACKNOWLEDGMENT

We would like to thank Professor J. H. Van Vleck for suggesting the subject and for valuable discussions of our experimental results.

<sup>4</sup> S. Velayos, *Anales Real Soc. Fis. Quim. (Madrid)* **33**, 5 (1935).